### The Story Continues... Systems with Continuous Diffraction

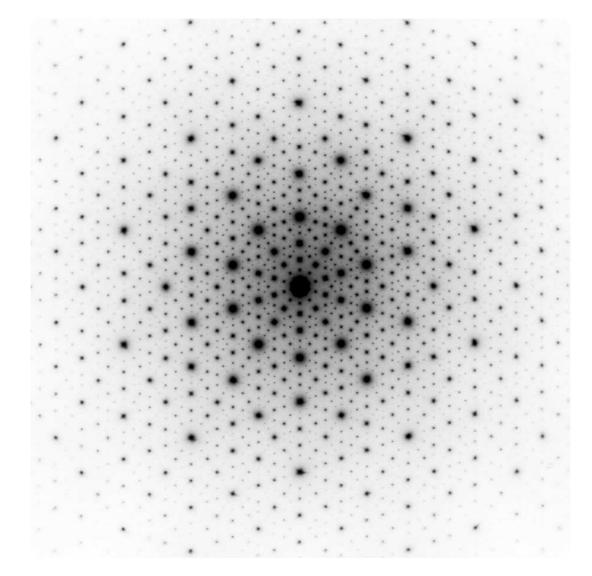
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(joint work with Michael Baake)

# Menue

- Diffraction theory
- Homometry
- Coin tossing
- Rudin-Shapiro
  - sequence
  - autocorrelation
  - diffraction
- Bernoullisation
- Outlook



### **Diffraction theory**

Setting:  $\omega \land \gamma = \omega \circledast \widetilde{\omega} \land \widehat{\gamma} \land \omega$ 

Dirac comb on  $\mathbb{Z}$ :

$$\omega = \sum_{n \in \mathbb{Z}} w(n) \, \delta_n \quad \curvearrowleft \quad \gamma = \sum_{m \in \mathbb{Z}} \eta(m) \, \delta_m$$

Autocorrelation coefficients:

$$\eta(m) = \lim_{N \to \infty} \frac{1}{2N+1} \sum_{n=-N}^{N} w(n) w(n+m)$$

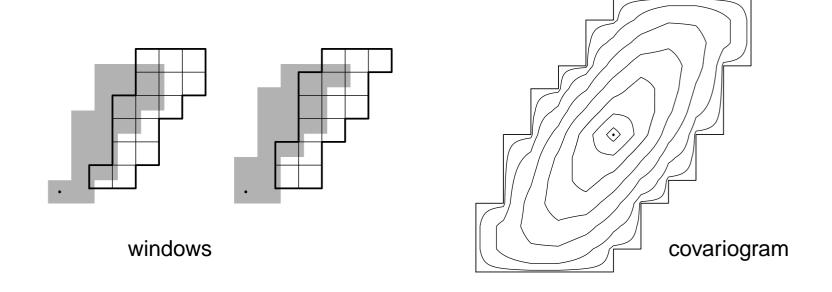
# Homometry

Problem: distinct structures with identical autocorrelation

Example 1:
$$\delta_{6\mathbb{Z}} * \sum_{j=0}^{5} c_j \, \delta_j$$
 $j$ 012345 $c_j$ 112542453114 $c_j$ 102139463517

same correlations up to order 5 (Grünbaum & Moore)

Example 2: homometric models sets with distinct windows



# **Coin tossing sequence**

Sequence: i.i.d. random variables  $W_n \in \{\pm 1\}$ with probabilities p and 1-p

**Metric entropy:**  $H(p) = -p \log(p) - (1-p) \log(1-p)$ 

Autocorrelation:  $\gamma_{\mathrm{B}} = \sum_{m \in \mathbb{Z}} \eta_{\mathrm{B}}(m) \delta_m$  with

$$\eta_{\rm B}(m) := \lim_{N \to \infty} \frac{1}{2N+1} \sum_{n=-N}^{N} W_n W_{n+m} \stackrel{\text{(a.s.)}}{=} \begin{cases} 1, & m = 0\\ (2p-1)^2, & m \neq 0 \end{cases}$$

(strong law of large numbers)

Diffraction measure:

$$\left| \begin{array}{c} \widehat{\gamma_{\mathrm{B}}} \stackrel{\text{(a.s.)}}{=} (2p-1)^2 \delta_{\mathbb{Z}} + 4p(1-p) \lambda \end{array} \right|$$

# **Rudin-Shapiro sequence**

Substitution:  $\varrho: a \mapsto ac, b \mapsto dc, c \mapsto ab, d \mapsto db$ 

Fixed point: 
$$b|a \xrightarrow{\varrho^2} dbab|acab \xrightarrow{\varrho^2} \dots \longrightarrow u = \varrho^2(u)$$

**Reduction:**  $\varphi: a, c \mapsto 1, b, d \mapsto -1, |w:=\varphi(u)|$ 

#### 

Alternative description: 
$$w(-1) = -1, w(0) = 1,$$
 with

$$w(4n+\ell) = \begin{cases} w(n), & \text{for } \ell \in \{0,1\}\\ (-1)^{n+\ell} w(n), & \text{for } \ell \in \{2,3\} \end{cases}$$

Autocorrelation:  $\gamma$ 

$$_{\mathrm{RS}} = \sum_{m \in \mathbb{Z}} \eta(m) \delta_m$$

### **Rudin-Shapiro autocorrelation**

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Define: 
$$\eta(m) \\ \vartheta(m) \end{cases}$$
 :=  $\lim_{N \to \infty} \frac{1}{2N+1} \sum_{n=-N}^{N} w(n) w(n+m) \begin{cases} 1 \\ (-1)^n \end{cases}$ 

(all limits exist by Birkhoff's ergodic theorem)

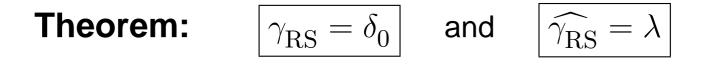
Recursion:  $\eta(0) = 1, \ \vartheta(0) = 0$ , and

$$\eta(4m) = \frac{1+(-1)^m}{2} \eta(m), \qquad \eta(4m+2) = 0,$$
  
$$\eta(4m+1) = \frac{1-(-1)^m}{4} \eta(m) + \frac{(-1)^m}{4} \vartheta(m) - \frac{1}{4} \vartheta(m+1),$$
  
$$\eta(4m+3) = \frac{1+(-1)^m}{4} \eta(m+1) - \frac{(-1)^m}{4} \vartheta(m) + \frac{1}{4} \vartheta(m+1),$$

$$\vartheta(4m) = 0, \qquad \vartheta(4m+2) = \frac{(-1)^m}{2} \vartheta(m) + \frac{1}{2} \vartheta(m+1),$$
$$\vartheta(4m+1) = \frac{1-(-1)^m}{4} \eta(m) - \frac{(-1)^m}{4} \vartheta(m) + \frac{1}{4} \vartheta(m+1),$$
$$\vartheta(4m+3) = -\frac{1+(-1)^m}{4} \eta(m+1) - \frac{(-1)^m}{4} \vartheta(m) + \frac{1}{4} \vartheta(m+1).$$

# **Rudin-Shapiro diffraction**

Unique solution:  $\vartheta(1) = 0 \quad \curvearrowright \quad \vartheta(m) = 0 \text{ for all } m \in \mathbb{Z}$ and  $\eta(m) = 0$  for all  $m \neq 0$ 



⇒ homometric with coin tossing for  $p = \frac{1}{2}$ , but zero entropy !

General weights:  $h_{\pm}$  instead of  $\pm 1$ :

$$\widehat{\gamma_h} = \left|\frac{h_+ + h_-}{2}\right|^2 \delta_{\mathbb{Z}} + \left|\frac{h_+ - h_-}{2}\right|^2 \lambda$$

### **Bernoullisation**

Sequence:  $S \in \{\pm 1\}^{\mathbb{Z}}$  (assumed ergodic) with Dirac comb  $\omega_S = \sum_{n \in \mathbb{Z}} S_n \, \delta_n$ and autocorrelation  $\gamma_S$ 

Bernoullisation:  $\omega := \sum_{n \in \mathbb{Z}} S_n W_n \, \delta_n \qquad (W_n \in \{\pm 1\})$ 

Autocorrelation: 
$$\gamma \stackrel{\text{(a.s.)}}{=} (2p-1)^2 \gamma_S + 4p(1-p) \delta_0$$

(strong law of large numbers)

Application: Rudin-Shapiro, with  $\gamma_S = \gamma_{RS} = \delta_0$  $\sim \gamma = \delta_0$  independently of p

 $\land$  diffraction  $|\widehat{\gamma} \equiv \lambda|$ 

 $\land$  homometric, irrespective of entropy

# Outlook

- Diffraction as useful tool
- Continuous spectra accessible
- Homometry more difficult
- Insensitivity to entropy
- Generalisation beyond lattice systems
- Extension to higher dimension
- Lower rank entropy (Ledrappier)
- Randomness with interaction

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