

Infinite Quantum Graphs

Aleksey Kostenko

University of Ljubljana, Slovenia
& University of Vienna, Austria

Differential Operators on Graphs and Waveguides
Graz, Austria

February 27, 2019



Univerza v Ljubljani



universität
wien

FWF

Der Wissenschaftsfonds.

Combinatorial and Metric Graphs

Definition

A (combinatorial) graph is the set of vertices \mathcal{V} and edges \mathcal{E} , $\mathcal{G}_d = (\mathcal{V}, \mathcal{E})$.

For $u, v \in \mathcal{V}$ we shall write $u \sim v$ if there is $e_{u,v} \in \mathcal{E}$ connecting u and v .
The function $\text{deg}: \mathcal{V} \rightarrow \mathbb{Z}_{\geq 1} \cup \{\infty\}$ defined by

$$\text{deg}: v \mapsto \#\{u \in \mathcal{V} \mid u \sim v\} = \#\mathcal{E}_v$$

is called **the (combinatorial) degree**, where $\mathcal{E}_v := \{e_{u,v} \in \mathcal{E} \mid u \sim v\}$.

Assumptions

Combinatorial and Metric Graphs

Definition

A (combinatorial) graph is the set of vertices \mathcal{V} and edges \mathcal{E} , $\mathcal{G}_d = (\mathcal{V}, \mathcal{E})$.

For $u, v \in \mathcal{V}$ we shall write $u \sim v$ if there is $e_{u,v} \in \mathcal{E}$ connecting u and v .
The function $\text{deg}: \mathcal{V} \rightarrow \mathbb{Z}_{\geq 1} \cup \{\infty\}$ defined by

$$\text{deg}: v \mapsto \#\{u \in \mathcal{V} \mid u \sim v\} = \#\mathcal{E}_v$$

is called **the (combinatorial) degree**, where $\mathcal{E}_v := \{e_{u,v} \in \mathcal{E} \mid u \sim v\}$.

Assumptions

- \mathcal{V} and \mathcal{E} are at most countable

Combinatorial and Metric Graphs

Definition

A (combinatorial) graph is the set of vertices \mathcal{V} and edges \mathcal{E} , $\mathcal{G}_d = (\mathcal{V}, \mathcal{E})$.

For $u, v \in \mathcal{V}$ we shall write $u \sim v$ if there is $e_{u,v} \in \mathcal{E}$ connecting u and v .
The function $\text{deg}: \mathcal{V} \rightarrow \mathbb{Z}_{\geq 1} \cup \{\infty\}$ defined by

$$\text{deg}: v \mapsto \#\{u \in \mathcal{V} \mid u \sim v\} = \#\mathcal{E}_v$$

is called **the (combinatorial) degree**, where $\mathcal{E}_v := \{e_{u,v} \in \mathcal{E} \mid u \sim v\}$.

Assumptions

- \mathcal{V} and \mathcal{E} are at most countable
- \mathcal{G}_d is connected and locally finite ($\text{deg}(v) < \infty$ for all $v \in \mathcal{V}$)

Combinatorial and Metric Graphs

Definition

A (combinatorial) graph is the set of vertices \mathcal{V} and edges \mathcal{E} , $\mathcal{G}_d = (\mathcal{V}, \mathcal{E})$.

For $u, v \in \mathcal{V}$ we shall write $u \sim v$ if there is $e_{u,v} \in \mathcal{E}$ connecting u and v .
The function $\deg: \mathcal{V} \rightarrow \mathbb{Z}_{\geq 1} \cup \{\infty\}$ defined by

$$\deg: v \mapsto \#\{u \in \mathcal{V} \mid u \sim v\} = \#\mathcal{E}_v$$

is called the (combinatorial) degree, where $\mathcal{E}_v := \{e_{u,v} \in \mathcal{E} \mid u \sim v\}$.

Assumptions

- \mathcal{V} and \mathcal{E} are at most countable
- \mathcal{G}_d is connected and locally finite ($\deg(v) < \infty$ for all $v \in \mathcal{V}$)
- No loops or multiple edges

Combinatorial and Metric Graphs

Definition

A (combinatorial) graph is the set of vertices \mathcal{V} and edges \mathcal{E} , $\mathcal{G}_d = (\mathcal{V}, \mathcal{E})$.

For $u, v \in \mathcal{V}$ we shall write $u \sim v$ if there is $e_{u,v} \in \mathcal{E}$ connecting u and v .
The function $\text{deg}: \mathcal{V} \rightarrow \mathbb{Z}_{\geq 1} \cup \{\infty\}$ defined by

$$\text{deg}: v \mapsto \#\{u \in \mathcal{V} \mid u \sim v\} = \#\mathcal{E}_v$$

is called **the (combinatorial) degree**, where $\mathcal{E}_v := \{e_{u,v} \in \mathcal{E} \mid u \sim v\}$.

Assumptions

- \mathcal{V} and \mathcal{E} are at most countable
- \mathcal{G}_d is connected and locally finite ($\text{deg}(v) < \infty$ for all $v \in \mathcal{V}$)
- No loops or multiple edges

Definition

If every edge $e \in \mathcal{E}$ is assigned with a length $|e| \in (0, \infty)$, then $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$ is called a **metric graph**

Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$, we can identify each edge $e \in \mathcal{E}$ with an interval $(0, |e|)$ and hence introduce the Hilbert space

$$L^2(\mathcal{G}) := \bigoplus_{e \in \mathcal{E}} L^2(e) = \left\{ f = \{f_e\}_{e \in \mathcal{E}} : f_e \in L^2(e), \sum_{e \in \mathcal{E}} \|f_e\|_{L^2(e)}^2 < \infty \right\}$$

Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$, we can identify each edge $e \in \mathcal{E}$ with an interval $(0, |e|)$ and hence introduce the Hilbert space

$$L^2(\mathcal{G}) := \bigoplus_{e \in \mathcal{E}} L^2(e) = \left\{ f = \{f_e\}_{e \in \mathcal{E}} : f_e \in L^2(e), \sum_{e \in \mathcal{E}} \|f_e\|_{L^2(e)}^2 < \infty \right\}$$

Next equip \mathcal{G} with a Schrödinger-type operator $\mathbf{H}_{\max} := \bigoplus_{e \in \mathcal{E}} H_e$, where:

$$H_e = -\frac{d^2}{dx_e^2}, \quad \text{dom}(H_e) = H^2(e).$$

Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$, we can identify each edge $e \in \mathcal{E}$ with an interval $(0, |e|)$ and hence introduce the Hilbert space

$$L^2(\mathcal{G}) := \bigoplus_{e \in \mathcal{E}} L^2(e) = \left\{ f = \{f_e\}_{e \in \mathcal{E}} : f_e \in L^2(e), \sum_{e \in \mathcal{E}} \|f_e\|_{L^2(e)}^2 < \infty \right\}$$

Next equip \mathcal{G} with a Schrödinger-type operator $\mathbf{H}_{\max} := \bigoplus_{e \in \mathcal{E}} H_e$, where:

$$H_e = -\frac{d^2}{dx_e^2} + V(x_e), \quad \text{dom}(H_e) = \mathcal{D}_{\max}(e).$$

Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$, we can identify each edge $e \in \mathcal{E}$ with an interval $(0, |e|)$ and hence introduce the Hilbert space

$$L^2(\mathcal{G}) := \bigoplus_{e \in \mathcal{E}} L^2(e) = \left\{ f = \{f_e\}_{e \in \mathcal{E}} : f_e \in L^2(e), \sum_{e \in \mathcal{E}} \|f_e\|_{L^2(e)}^2 < \infty \right\}$$

Next equip \mathcal{G} with a Schrödinger-type operator $\mathbf{H}_{\max} := \bigoplus_{e \in \mathcal{E}} \mathbf{H}_e$, where:

$$\mathbf{H}_e = \left(\frac{1}{i} \frac{d}{dx_e} - A(x_e) \right)^2 + V(x_e), \quad \text{dom}(\mathbf{H}_e) = \mathcal{D}_{\max}(e).$$

Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$, we can identify each edge $e \in \mathcal{E}$ with an interval $(0, |e|)$ and hence introduce the *weighted* Hilbert space

$$L^2(\mathcal{G}; \mu) := \bigoplus_{e \in \mathcal{E}} L^2(e; \mu_e)$$

Next equip \mathcal{G} with a Schrödinger-type operator $\mathbf{H}_{\max} := \bigoplus_{e \in \mathcal{E}} \mathbf{H}_e$, where:

$$\mathbf{H}_e = -\frac{1}{\mu_e} \frac{d}{dx_e} \nu_e \frac{d}{dx_e}, \quad \text{dom}(\mathbf{H}_e) = H^2(e).$$

Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$, we can identify each edge $e \in \mathcal{E}$ with an interval $(0, |e|)$ and hence introduce the Hilbert space

$$L^2(\mathcal{G}) := \bigoplus_{e \in \mathcal{E}} L^2(e) = \left\{ f = \{f_e\}_{e \in \mathcal{E}} \mid f_e \in L^2(e), \sum_{e \in \mathcal{E}} \|f_e\|_{L^2(e)}^2 < \infty \right\}$$

Next equip \mathcal{G} with a Schrödinger-type operator $\mathbf{H}_{\max} := \bigoplus_{e \in \mathcal{E}} H_e$, where:

$$H_e = -\frac{d^2}{dx_e^2}, \quad \text{dom}(H_e) = H^2(e).$$

Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$, we can identify each edge $e \in \mathcal{E}$ with an interval $(0, |e|)$ and hence introduce the Hilbert space

$$L^2(\mathcal{G}) := \bigoplus_{e \in \mathcal{E}} L^2(e) = \left\{ f = \{f_e\}_{e \in \mathcal{E}} \mid f_e \in L^2(e), \sum_{e \in \mathcal{E}} \|f_e\|_{L^2(e)}^2 < \infty \right\}$$

Next equip \mathcal{G} with a Schrödinger-type operator $\mathbf{H}_{\max} := \bigoplus_{e \in \mathcal{E}} \mathbf{H}_e$, where:

$$\mathbf{H}_e = -\frac{d^2}{dx_e^2}, \quad \text{dom}(\mathbf{H}_e) = H^2(e).$$

To give \mathbf{H} the meaning of a quantum mechanical energy operator, it must be self-adjoint, that is, we need to add boundary conditions at the vertices.

Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$, we can identify each edge $e \in \mathcal{E}$ with an interval $(0, |e|)$ and hence introduce the Hilbert space

$$L^2(\mathcal{G}) := \bigoplus_{e \in \mathcal{E}} L^2(e) = \left\{ f = \{f_e\}_{e \in \mathcal{E}} \mid f_e \in L^2(e), \sum_{e \in \mathcal{E}} \|f_e\|_{L^2(e)}^2 < \infty \right\}$$

Next equip \mathcal{G} with a Schrödinger-type operator $\mathbf{H}_{\max} := \bigoplus_{e \in \mathcal{E}} H_e$, where:

$$H_e = -\frac{d^2}{dx_e^2}, \quad \text{dom}(H_e) = H^2(e).$$

To give \mathbf{H} the meaning of a quantum mechanical energy operator, it must be self-adjoint, that is, we need to add boundary conditions at the vertices.

$$f_e(v) := \lim_{x_e \rightarrow v} f(x_e), \quad f'_e(v) := \lim_{x_e \rightarrow v} \frac{f(x_e) - f_e(v)}{|x_e - v|}.$$

are well defined for all $f \in \text{dom}(\mathbf{H}_{\max})$.

Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$, we can identify each edge $e \in \mathcal{E}$ with an interval $(0, |e|)$ and hence introduce the Hilbert space

$$L^2(\mathcal{G}) := \bigoplus_{e \in \mathcal{E}} L^2(e) = \left\{ f = \{f_e\}_{e \in \mathcal{E}} \mid f_e \in L^2(e), \sum_{e \in \mathcal{E}} \|f_e\|_{L^2(e)}^2 < \infty \right\}$$

Next equip \mathcal{G} with a Schrödinger-type operator $\mathbf{H}_{\max} := \bigoplus_{e \in \mathcal{E}} H_e$, where:

$$H_e = -\frac{d^2}{dx_e^2}, \quad \text{dom}(H_e) = H^2(e).$$

To give \mathbf{H} the meaning of a quantum mechanical energy operator, it must be self-adjoint, that is, we need to add boundary conditions at the vertices.

(standard) Kirchhoff conditions: For all $v \in \mathcal{V}$

$$\begin{cases} f \text{ is continuous at } v, \\ \sum_{e \in \mathcal{E}_v} f'_e(v) = 0. \end{cases}$$

Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$, we can identify each edge $e \in \mathcal{E}$ with an interval $(0, |e|)$ and hence introduce the Hilbert space

$$L^2(\mathcal{G}) := \bigoplus_{e \in \mathcal{E}} L^2(e) = \left\{ f = \{f_e\}_{e \in \mathcal{E}} \mid f_e \in L^2(e), \sum_{e \in \mathcal{E}} \|f_e\|_{L^2(e)}^2 < \infty \right\}$$

Next equip \mathcal{G} with a Schrödinger-type operator $\mathbf{H}_{\max} := \bigoplus_{e \in \mathcal{E}} H_e$, where:

$$H_e = -\frac{d^2}{dx_e^2}, \quad \text{dom}(H_e) = H^2(e).$$

To give \mathbf{H} the meaning of a quantum mechanical energy operator, it must be self-adjoint, that is, we need to add boundary conditions at the vertices.

Definition

A **quantum graph** is a metric graph equipped with the operator \mathbf{H} acting as the negative second order derivative along edges and accompanied by Kirchhoff vertex conditions

Equilateral Quantum Graphs

Suppose $|e| = 1$ for all $e \in \mathcal{E}$.

Equilateral Quantum Graphs

Suppose $|e| = 1$ for all $e \in \mathcal{E}$.

Then the Kirchhoff Laplacian $\mathbf{H}_{\text{equil}}$ is self-adjoint.

Equilateral Quantum Graphs

Suppose $|e| = 1$ for all $e \in \mathcal{E}$.

Then the Kirchhoff Laplacian $\mathbf{H}_{\text{equil}}$ is self-adjoint.

Problem

Spectral analysis of $\mathbf{H}_{\text{equil}}$?

Equilateral Quantum Graphs

Suppose $|e| = 1$ for all $e \in \mathcal{E}$.

Then the Kirchhoff Laplacian $\mathbf{H}_{\text{equil}}$ is self-adjoint.

Problem

Spectral analysis of $\mathbf{H}_{\text{equil}}$?

Define the **normalized/physical Laplacian** on \mathcal{G}_d by

$$(\tau_{\text{norm}} f)(v) := \frac{1}{\text{deg}(v)} \sum_{u \sim v} f(v) - f(u), \quad v \in \mathcal{V}.$$

τ_{norm} generates a **bounded self-adjoint operator** h_{norm} in $\ell^2(\mathcal{V}; \text{deg})$.



R. Courant, K. Friedrichs and H. Lewy, *Über die partiellen Differenzengleichungen der mathematischen Physik*, Math. Ann. (1928)

Equilateral Quantum Graphs

Suppose $|e| = 1$ for all $e \in \mathcal{E}$.

Then the Kirchhoff Laplacian $\mathbf{H}_{\text{equil}}$ is self-adjoint.

Problem

Spectral analysis of $\mathbf{H}_{\text{equil}}$?

Define the **normalized/physical Laplacian** on \mathcal{G}_d by

$$(\tau_{\text{norm}} f)(v) := \frac{1}{\deg(v)} \sum_{u \sim v} f(v) - f(u), \quad v \in \mathcal{V}.$$

τ_{norm} generates a **bounded self-adjoint operator** h_{norm} in $\ell^2(\mathcal{V}; \deg)$.

-  R. Courant, K. Friedrichs and H. Lewy, *Über die partiellen Differenzengleichungen der mathematischen Physik*, Math. Ann. (1928)
-  Y. Colin de Verdière, *Spectres de Graphes*, SMF, Paris, 1998.
-  W. Woess, *Random Walks on Infinite Graphs and Groups*, CUP, 2000.

Equilateral Quantum Graphs

Suppose $|e| = 1$ for all $e \in \mathcal{E}$.

Then the Kirchhoff Laplacian $\mathbf{H}_{\text{equil}}$ is self-adjoint.

Problem

Spectral analysis of $\mathbf{H}_{\text{equil}}$?

Define the **normalized/physical Laplacian** on \mathcal{G}_d by

$$(\tau_{\text{norm}} f)(v) := \frac{1}{\deg(v)} \sum_{u \sim v} f(v) - f(u), \quad v \in \mathcal{V}.$$

τ_{norm} generates a **bounded self-adjoint operator** h_{norm} in $\ell^2(\mathcal{V}; \deg)$.

Theorem (von Below'87,..., Cattaneo, Exner,..., Pankrashkin'2012)

$$\sigma_j(\mathbf{H}_{\text{equil}}) \setminus \sigma_D = \{\lambda \notin \sigma_D \mid 1 - \cos(\sqrt{\lambda}) \in \sigma_j(h_{\text{norm}})\}, \quad j \in \{\text{p, ess, ac, sc}\}$$

with $\sigma_D = \{(\pi n)^2\}_{n \in \mathbb{N}}$.

Equilateral Quantum Graphs

Suppose $|e| = 1$ for all $e \in \mathcal{E}$.

Then the Kirchhoff Laplacian $\mathbf{H}_{\text{equil}}$ is self-adjoint.

Problem

Spectral analysis of $\mathbf{H}_{\text{equil}}$?

Define the **normalized/physical Laplacian** on \mathcal{G}_d by

$$(\tau_{\text{norm}} f)(v) := \frac{1}{\deg(v)} \sum_{u \sim v} f(v) - f(u), \quad v \in \mathcal{V}.$$

τ_{norm} generates a **bounded self-adjoint operator** h_{norm} in $\ell^2(\mathcal{V}; \deg)$.

Theorem (von Below'87,..., Cattaneo, Exner,..., Pankrashkin'2012)

$$\sigma_j(\mathbf{H}_{\text{equil}}) \setminus \sigma_D = \{\lambda \notin \sigma_D \mid 1 - \cos(\sqrt{\lambda}) \in \sigma_j(h_{\text{norm}})\}, \quad j \in \{p, \text{ess}, \text{ac}, \text{sc}\}$$

with $\sigma_D = \{(\pi n)^2\}_{n \in \mathbb{N}}$. $\mathbf{H}_{\text{equil}}$ and h_{norm} are “locally” unitarily equivalent.

(Non-equilateral) Quantum Graphs

Problem

Does there exist an analogous statement for **non-equilateral** graphs?

(Non-equilateral) Quantum Graphs

Problem

Does there exist an analogous statement for **non-equilateral** graphs?

Consider the (minimal) discrete Laplacian h_G defined on $\ell^2(\mathcal{V}; m)$ by

$$(\tau_G f)(v) := \frac{1}{m(v)} \sum_{u \sim v} \frac{f(v) - f(u)}{|e_{u,v}|}, \quad m(v) = \sum_{e \in \mathcal{E}_v} |e|.$$

τ_G is the normalized Laplacian **iff** $|e| = 1$ for all $e \in \mathcal{E}$ and $\alpha \equiv 0$.

(Non-equilateral) Quantum Graphs

Problem

Does there exist an analogous statement for **non-equilateral** graphs?

Consider the (minimal) discrete Laplacian h_G defined on $\ell^2(\mathcal{V}; m)$ by

$$(\tau_G f)(v) := \frac{1}{m(v)} \sum_{u \sim v} \frac{f(v) - f(u)}{|e_{u,v}|}, \quad m(v) = \sum_{e \in \mathcal{E}_v} |e|.$$

τ_G is the normalized Laplacian **iff** $|e| = 1$ for all $e \in \mathcal{E}$ and $\alpha \equiv 0$.

Theorem (E. B. Davies'1992)

h_G is bounded \Leftrightarrow the weighted degree **Deg** is bounded on \mathcal{V} ,

$$\text{Deg}: v \mapsto \frac{1}{m(v)} \sum_{u \sim v} \frac{1}{|e_{u,v}|} = \frac{\sum_{e \in \mathcal{E}_v} 1/|e|}{\sum_{e \in \mathcal{E}_v} |e|}$$

(Non-equilateral) Quantum Graphs

Problem

Does there exist an analogous statement for **non-equilateral** graphs?

Consider the (minimal) discrete Laplacian h_G defined on $\ell^2(\mathcal{V}; m)$ by

$$(\tau_G f)(v) := \frac{1}{m(v)} \sum_{u \sim v} \frac{f(v) - f(u)}{|e_{u,v}|}, \quad m(v) = \sum_{e \in \mathcal{E}_v} |e|.$$

τ_G is the normalized Laplacian **iff** $|e| = 1$ for all $e \in \mathcal{E}$ and $\alpha \equiv 0$.

Theorem (E. B. Davies'1992)

h_G is bounded \Leftrightarrow the weighted degree **Deg** is bounded on \mathcal{V} ,

$$\text{Deg}: v \mapsto \frac{1}{m(v)} \sum_{u \sim v} \frac{1}{|e_{u,v}|} = \frac{\sum_{e \in \mathcal{E}_v} 1/|e|}{\sum_{e \in \mathcal{E}_v} |e|}$$

Note that **Deg** is bounded on \mathcal{V} if $\underline{l}_*(\mathcal{E}) := \inf_{e \in \mathcal{E}} |e| > 0$.

(Non-equilateral) Quantum Graphs

Theorem 1 (Exner, AK, Malamud, Neidhardt'2018)

Let \mathcal{G} be a metric graph with $\ell^*(\mathcal{G}) := \sup_{e \in \mathcal{E}} |e| < \infty$. Then:

(i) $\mathbf{H}_{\mathcal{G}}$ is self-adjoint $\iff h_{\mathcal{G}}$ is self-adjoint,

(Non-equilateral) Quantum Graphs

Theorem 1 (Exner, AK, Malamud, Neidhardt'2018)

Let \mathcal{G} be a metric graph with $\ell^*(\mathcal{G}) := \sup_{e \in \mathcal{E}} |e| < \infty$. Then:

- (i) $\mathbf{H}_{\mathcal{G}}$ is self-adjoint $\iff h_{\mathcal{G}}$ is self-adjoint,
- (ii) $\inf \sigma(\mathbf{H}_{\mathcal{G}}) > 0 \iff \inf \sigma(h_{\mathcal{G}}) > 0$.
- (iii) $\inf \sigma_{\text{ess}}(\mathbf{H}_{\mathcal{G}}) > 0 \iff \inf \sigma_{\text{ess}}(h_{\mathcal{G}}) > 0$.

(Non-equilateral) Quantum Graphs

Theorem 1 (Exner, AK, Malamud, Neidhardt'2018)

Let \mathcal{G} be a metric graph with $\ell^*(\mathcal{G}) := \sup_{e \in \mathcal{E}} |e| < \infty$. Then:

- (i) $\mathbf{H}_{\mathcal{G}}$ is self-adjoint $\iff h_{\mathcal{G}}$ is self-adjoint,
- (ii) $\inf \sigma(\mathbf{H}_{\mathcal{G}}) > 0 \iff \inf \sigma(h_{\mathcal{G}}) > 0$.
- (iii) $\inf \sigma_{\text{ess}}(\mathbf{H}_{\mathcal{G}}) > 0 \iff \inf \sigma_{\text{ess}}(h_{\mathcal{G}}) > 0$.
- (iv) $\sigma(\mathbf{H}_{\mathcal{G}})$ is discrete $\iff \sigma(h_{\mathcal{G}})$ is discrete and $\#\{e \in \mathcal{E} : |e| > \varepsilon\}$ is finite for all $\varepsilon > 0$.

(Non-equilateral) Quantum Graphs

Theorem 1 (Exner, AK, Malamud, Neidhardt'2018)

Let \mathcal{G} be a metric graph with $\ell^*(\mathcal{G}) := \sup_{e \in \mathcal{E}} |e| < \infty$. Then:

- (i) $\mathbf{H}_{\mathcal{G}}$ is self-adjoint $\iff h_{\mathcal{G}}$ is self-adjoint,
- (ii) $\inf \sigma(\mathbf{H}_{\mathcal{G}}) > 0 \iff \inf \sigma(h_{\mathcal{G}}) > 0$.
- (iii) $\inf \sigma_{\text{ess}}(\mathbf{H}_{\mathcal{G}}) > 0 \iff \inf \sigma_{\text{ess}}(h_{\mathcal{G}}) > 0$.
- (iv) $\sigma(\mathbf{H}_{\mathcal{G}})$ is discrete $\iff \sigma(h_{\mathcal{G}})$ is discrete and $\#\{e \in \mathcal{E} : |e| > \varepsilon\}$ is finite for all $\varepsilon > 0$.
- (v)

$$\|e^{-t h_{\mathcal{G}}}\|_{\ell^1 \rightarrow \ell^\infty} \leq C_1 t^{-D/2}, \quad t > 0,$$

for some $D > 2$ if and only if

$$\|e^{-t \mathbf{H}_{\mathcal{G}}}\|_{L^1 \rightarrow L^\infty} \leq C_2 t^{-D/2}, \quad t > 0.$$



P. Exner, A. Kostenko, M. Malamud, & H. Neidhardt, *Spectral theory of infinite quantum graphs*, Ann. Henri Poincaré **19**, no. 11, (2018).

(Non-equilateral) Quantum Graphs

Theorem 1 (Exner, AK, Malamud, Neidhardt'2018)

Let \mathcal{G} be a metric graph with $\ell^*(\mathcal{G}) := \sup_{e \in \mathcal{E}} |e| < \infty$. Then:

- (i) $\mathbf{H}_{\mathcal{G}}$ is self-adjoint $\iff h_{\mathcal{G}}$ is self-adjoint,
- (ii) $\inf \sigma(\mathbf{H}_{\mathcal{G}}) > 0 \iff \inf \sigma(h_{\mathcal{G}}) > 0$.
- (iii) $\inf \sigma_{\text{ess}}(\mathbf{H}_{\mathcal{G}}) > 0 \iff \inf \sigma_{\text{ess}}(h_{\mathcal{G}}) > 0$.
- (iv) $\sigma(\mathbf{H}_{\mathcal{G}})$ is discrete $\iff \sigma(h_{\mathcal{G}})$ is discrete and $\#\{e \in \mathcal{E} : |e| > \varepsilon\}$ is finite for all $\varepsilon > 0$.
- (v)

$$\|e^{-t h_{\mathcal{G}}}\|_{\ell^1 \rightarrow \ell^\infty} \leq C_1 t^{-D/2}, \quad t > 0,$$

for some $D > 2$ if and only if

$$\|e^{-t \mathbf{H}_{\mathcal{G}}}\|_{L^1 \rightarrow L^\infty} \leq C_2 t^{-D/2}, \quad t > 0.$$



A. Kostenko and N. Nicolussi, *Spectral estimates for infinite quantum graphs*, Calc. Var. Partial Differential Equations **58**, no. 1, (2019).

Quantum Graphs: Self-adjointness

For $p: \mathcal{E} \rightarrow (0, \infty)$, define a **path metric** ϱ_p on \mathcal{V} w.r.t. \mathcal{G} by

$$\varrho_p(u, v) := \inf_{\mathcal{P}=\{v_0, \dots, v_n\}: u=v_0, v=v_n} \sum_k p(e_{v_{k-1}, v_k}).$$

The infimum is taken over all paths connecting u and v .

Quantum Graphs: Self-adjointness

For $p: \mathcal{E} \rightarrow (0, \infty)$, define a **path metric** ϱ_p on \mathcal{V} w.r.t. \mathcal{G} by

$$\varrho_p(u, v) := \inf_{\mathcal{P}=\{v_0, \dots, v_n\}: u=v_0, v=v_n} \sum_k p(e_{v_{k-1}, v_k}).$$

The infimum is taken over all paths connecting u and v .

Examples

- Natural path metric ϱ_0 with $p_0: e \mapsto |e|$.

Quantum Graphs: Self-adjointness

For $p: \mathcal{E} \rightarrow (0, \infty)$, define a **path metric** ϱ_p on \mathcal{V} w.r.t. \mathcal{G} by

$$\varrho_p(u, v) := \inf_{\mathcal{P}=\{v_0, \dots, v_n\}: u=v_0, v=v_n} \sum_k p(e_{v_{k-1}, v_k}).$$

The infimum is taken over all paths connecting u and v .

Examples

- Natural path metric ϱ_0 with $p_0: e \mapsto |e|$.
- Star metric ϱ_m with $p_m: e_{u,v} \mapsto m(u) + m(v)$ with $m(v) = \sum_{e \in \mathcal{E}_v} |e|$

Quantum Graphs: Self-adjointness

For $p: \mathcal{E} \rightarrow (0, \infty)$, define a **path metric** ϱ_p on \mathcal{V} w.r.t. \mathcal{G} by

$$\varrho_p(u, v) := \inf_{\mathcal{P}=\{v_0, \dots, v_n\}: u=v_0, v=v_n} \sum_k p(e_{v_{k-1}, v_k}).$$

The infimum is taken over all paths connecting u and v .

Examples

- Natural path metric ϱ_0 with $p_0: e \mapsto |e|$.
- Star metric ϱ_m with $p_m: e_{u,v} \mapsto m(u) + m(v)$ with $m(v) = \sum_{e \in \mathcal{E}_v} |e|$

Hopf–Rinow-type Theorem

(\mathcal{V}, ϱ_p) is complete as a metric space \iff

(\mathcal{V}, ϱ_p) is geodesically complete \iff

The distance balls in (\mathcal{V}, ϱ_p) are finite (“*finite ball condition*”).



X. Huang, M. Keller, J. Masamune, R. Wojciechowski, *A note on self-adjoint extensions of the Laplacian on weighted graphs*, J. Funct. Anal. **265** (2013).

Quantum Graphs: Self-adjointness

For $p: \mathcal{E} \rightarrow (0, \infty)$, define a **path metric** ϱ_p on \mathcal{V} w.r.t. \mathcal{G} by

$$\varrho_p(u, v) := \inf_{\mathcal{P}=\{v_0, \dots, v_n\}: u=v_0, v=v_n} \sum_k p(e_{v_{k-1}, v_k}).$$

The infimum is taken over all paths connecting u and v .

Examples

- Natural path metric ϱ_0 with $p_0: e \mapsto |e|$.
- Star metric ϱ_m with $p_m: e_{u,v} \mapsto m(u) + m(v)$ with $m(v) = \sum_{e \in \mathcal{E}_v} |e|$

Theorem 2 (Exner–AK–Malamud–Neidhardt)

If (\mathcal{V}, ϱ_m) is complete as a metric space, then \mathbf{H} is self-adjoint.

Quantum Graphs: Self-adjointness

For $p: \mathcal{E} \rightarrow (0, \infty)$, define a **path metric** ϱ_p on \mathcal{V} w.r.t. \mathcal{G} by

$$\varrho_p(u, v) := \inf_{\mathcal{P}=\{v_0, \dots, v_n\}: u=v_0, v=v_n} \sum_k p(e_{v_{k-1}, v_k}).$$

The infimum is taken over all paths connecting u and v .

Examples

- Natural path metric ϱ_0 with $p_0: e \mapsto |e|$.
- Star metric ϱ_m with $p_m: e_{u,v} \mapsto m(u) + m(v)$ with $m(v) = \sum_{e \in \mathcal{E}_v} |e|$

Theorem 2 (Exner–AK–Malamud–Neidhardt)

If (\mathcal{V}, ϱ_m) is complete as a metric space, then \mathbf{H} is self-adjoint.

In particular, \mathbf{H} is self-adjoint if $\inf_{v \in \mathcal{V}} m(v) = \inf_{v \in \mathcal{V}} \sum_{e \in \mathcal{E}_v} |e| > 0$.



M. Keller and D. Lenz, *Dirichlet forms and stochastic completeness of graphs and subgraphs*, J. reine angew. Math. **666** (2012).

Quantum Graphs: Self-adjointness

For $p: \mathcal{E} \rightarrow (0, \infty)$, define a **path metric** ϱ_p on \mathcal{V} w.r.t. \mathcal{G} by

$$\varrho_p(u, v) := \inf_{\mathcal{P}=\{v_0, \dots, v_n\}: u=v_0, v=v_n} \sum_k p(e_{v_{k-1}, v_k}).$$

The infimum is taken over all paths connecting u and v .

Examples

- Natural path metric ϱ_0 with $p_0: e \mapsto |e|$.
- Star metric ϱ_m with $p_m: e_{u,v} \mapsto m(u) + m(v)$ with $m(v) = \sum_{e \in \mathcal{E}_v} |e|$

Theorem 2 (Exner–AK–Malamud–Neidhardt)

If (\mathcal{V}, ϱ_m) is complete as a metric space, then \mathbf{H} is self-adjoint.

In particular, \mathbf{H} is self-adjoint if $\inf_{v \in \mathcal{V}} m(v) = \inf_{v \in \mathcal{V}} \sum_{e \in \mathcal{E}_v} |e| > 0$.

Gaffney-type Theorem: (\mathcal{G}, ϱ_0) is complete $\Rightarrow \mathbf{H}_{\mathcal{G}}$ is self-adjoint.

Quantum Graphs: Self-adjointness

For $p: \mathcal{E} \rightarrow (0, \infty)$, define a **path metric** ϱ_p on \mathcal{V} w.r.t. \mathcal{G} by

$$\varrho_p(u, v) := \inf_{\mathcal{P}=\{v_0, \dots, v_n\}: u=v_0, v=v_n} \sum_k p(e_{v_{k-1}, v_k}).$$

The infimum is taken over all paths connecting u and v .

Examples

- Natural path metric ϱ_0 with $p_0: e \mapsto |e|$.
- Star metric ϱ_m with $p_m: e_{u,v} \mapsto m(u) + m(v)$ with $m(v) = \sum_{e \in \mathcal{E}_v} |e|$

Theorem 2 (Exner–AK–Malamud–Neidhardt)

If (\mathcal{V}, ϱ_m) is complete as a metric space, then \mathbf{H} is self-adjoint.

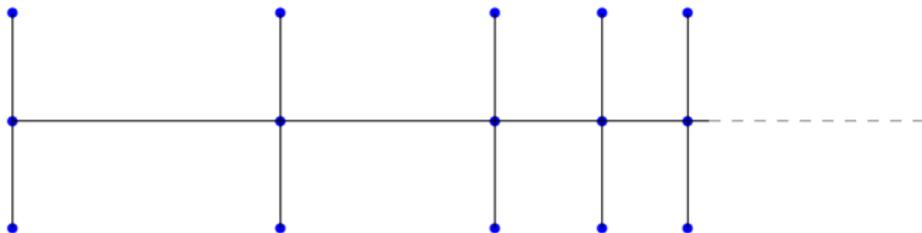
In particular, \mathbf{H} is self-adjoint if $\inf_{v \in \mathcal{V}} m(v) = \inf_{v \in \mathcal{V}} \sum_{e \in \mathcal{E}_v} |e| > 0$.

Gaffney-type Theorem: (\mathcal{G}, ϱ_0) is complete $\Rightarrow \mathbf{H}_{\mathcal{G}}$ is self-adjoint.

The standard assumption for infinite QG is $\inf_{e \in \mathcal{E}} |e| > 0!$

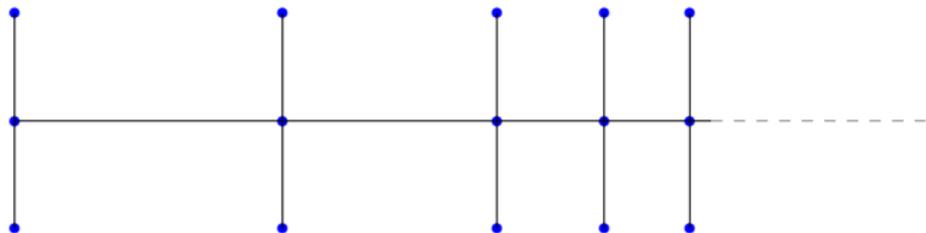
Quantum Graphs: Self-adjointness

Example 1.



Quantum Graphs: Self-adjointness

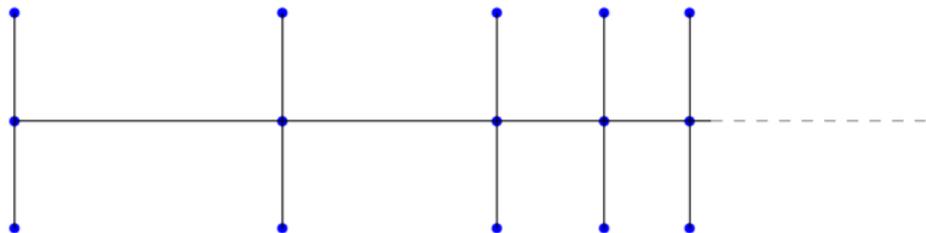
Example 1.



In Example 1, (\mathcal{V}, ϱ_m) is complete $\Leftrightarrow m(\mathcal{V}) = 2\text{vol}(\mathcal{G}) = 2 \sum_{e \in \mathcal{E}} |e| = \infty$.

Quantum Graphs: Self-adjointness

Example 1.



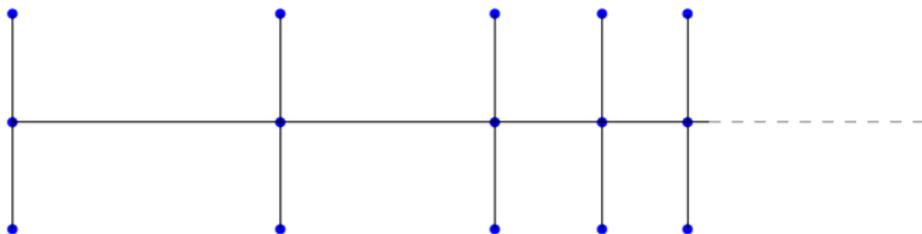
In Example 1, (\mathcal{V}, ϱ_m) is complete $\Leftrightarrow m(\mathcal{V}) = 2\text{vol}(\mathcal{G}) = 2 \sum_{e \in \mathcal{E}} |e| = \infty$.

Lemma

If $\text{vol}(\mathcal{G}) < \infty$, then \mathbf{H} is non-self-adjoint.

Quantum Graphs: Self-adjointness

Example 1.



In Example 1, (\mathcal{V}, ϱ_m) is complete $\Leftrightarrow m(\mathcal{V}) = 2\text{vol}(\mathcal{G}) = 2 \sum_{e \in \mathcal{E}} |e| = \infty$.

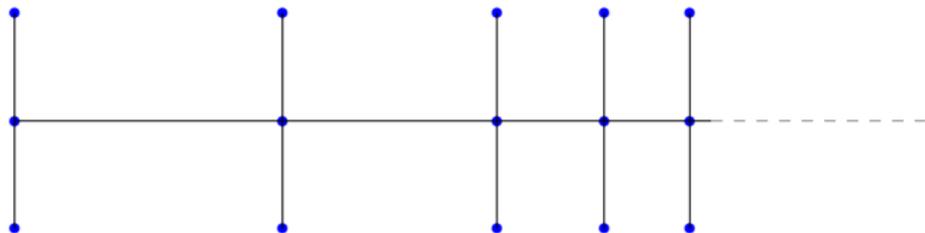
Lemma

If $\text{vol}(\mathcal{G}) < \infty$, then \mathbf{H} is non-self-adjoint.

Hence, in Example 1, \mathbf{H} is self-adjoint $\Leftrightarrow (\mathcal{V}, \varrho_m)$ is complete!

Quantum Graphs: Self-adjointness

Example 1.



In Example 1, (\mathcal{V}, ϱ_m) is complete $\Leftrightarrow m(\mathcal{V}) = 2\text{vol}(\mathcal{G}) = 2 \sum_{e \in \mathcal{E}} |e| = \infty$.

Lemma

If $\text{vol}(\mathcal{G}) < \infty$, then \mathbf{H} is non-self-adjoint.

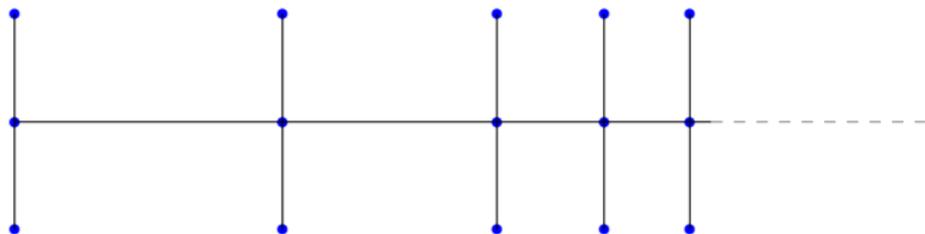
Hence, in Example 1, \mathbf{H} is self-adjoint $\Leftrightarrow (\mathcal{V}, \varrho_m)$ is complete!

Remark

The converse to Theorem 2 is not true!

Quantum Graphs: Self-adjointness

Example 1.



In Example 1, (\mathcal{V}, ϱ_m) is complete $\Leftrightarrow m(\mathcal{V}) = 2\text{vol}(\mathcal{G}) = 2 \sum_{e \in \mathcal{E}} |e| = \infty$.

Lemma

If $\text{vol}(\mathcal{G}) < \infty$, then \mathbf{H} is non-self-adjoint.

Hence, in Example 1, \mathbf{H} is self-adjoint $\Leftrightarrow (\mathcal{V}, \varrho_m)$ is complete!

Remark

The converse to Theorem 2 is not true!

For radially symmetric trees and antitrees, \mathbf{H} is self-adjoint $\Leftrightarrow m(\mathcal{V}) = \infty$.

Examples: Radially symmetric antitrees

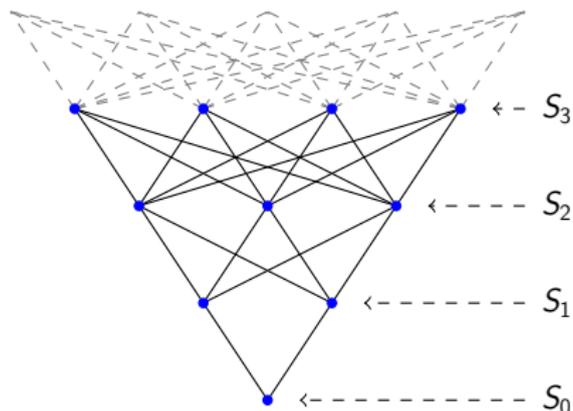


Figure: Example of an antitree \mathcal{A} with $s_n = n + 1$.

S_n is the n -th combinatorial sphere, and $s_n := \#S_n$.

\mathcal{A} is radially symmetric if edges connecting S_n with S_{n+1} have the same length, say ℓ_n , for all $n \geq 0$.

Examples: Radially symmetric antitrees

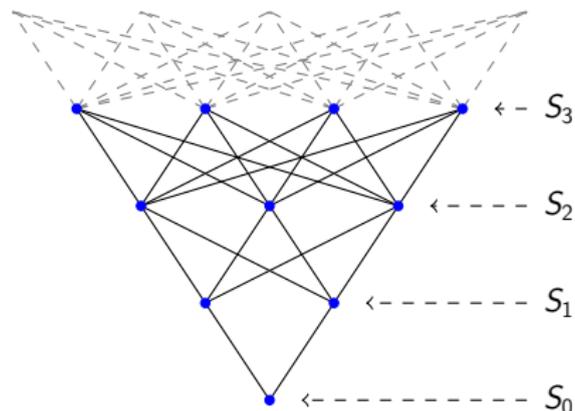


Figure: Example of an antitree \mathcal{A} with $s_n = n + 1$.

Theorem (AK–Nicolussi)

\mathbf{H} is self-adjoint $\iff \text{vol}(\mathcal{A}) = \sum_{n \geq 0} s_n s_{n+1} \ell_n = \infty$

Examples: Radially symmetric antitrees

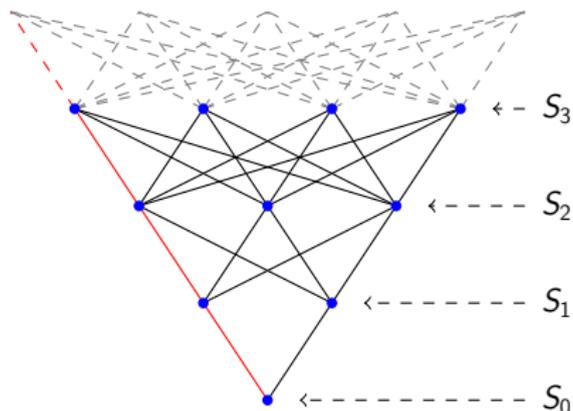


Figure: Example of an antitree \mathcal{A} with $s_n = n + 1$.

Theorem (AK–Nicolussi)

\mathbf{H} is self-adjoint $\iff \text{vol}(\mathcal{A}) = \sum_{n \geq 0} s_n s_{n+1} \ell_n = \infty$

(\mathcal{A}, ϱ_0) is complete $\iff \sum_{n \geq 0} \ell_n = \infty$.

Examples: Radially symmetric antitrees

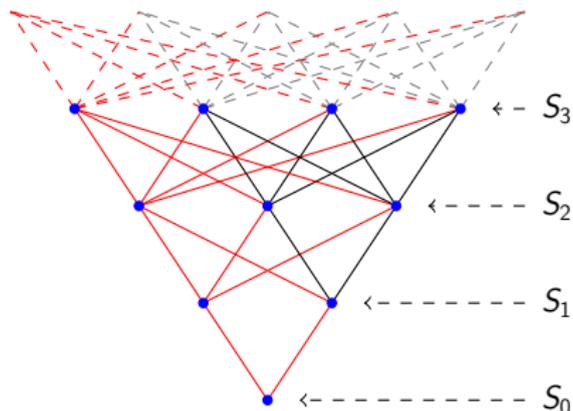


Figure: Example of an antitree \mathcal{A} with $s_n = n + 1$.

Theorem (AK–Nicolussi)

\mathbf{H} is self-adjoint $\iff \text{vol}(\mathcal{A}) = \sum_{n \geq 0} s_n s_{n+1} \ell_n = \infty$

(\mathcal{A}, ϱ_0) is complete $\iff \sum_{n \geq 0} \ell_n = \infty$.

(\mathcal{V}, ϱ_m) is complete $\iff \sum_{n \geq 0} (s_n + s_{n+1}) \ell_n = \infty$

Summary

- (i) \mathbf{H} is self-adjoint if (\mathcal{V}, ϱ_m) is complete.
- (ii) \mathbf{H} is non-self-adjoint if $\text{vol}(\mathcal{G}) = \sum_{e \in \mathcal{E}} |e| < \infty$.

Summary

- (i) \mathbf{H} is self-adjoint if (\mathcal{V}, ϱ_m) is complete.
- (ii) \mathbf{H} is non-self-adjoint if $\text{vol}(\mathcal{G}) = \sum_{e \in \mathcal{E}} |e| < \infty$.

Problems

- (i) Characterize metric graphs such that completeness of (\mathcal{V}, ϱ_m) is also necessary for self-adjointness.

Summary

- (i) \mathbf{H} is self-adjoint if (\mathcal{V}, ϱ_m) is complete.
- (ii) \mathbf{H} is non-self-adjoint if $\text{vol}(\mathcal{G}) = \sum_{e \in \mathcal{E}} |e| < \infty$.

Problems

- (i) Characterize metric graphs such that completeness of (\mathcal{V}, ϱ_m) is also necessary for self-adjointness.
- (ii) Characterize metric graphs such that $\text{vol}(\mathcal{G}) = \sum_{e \in \mathcal{E}} |e| = \infty$ is also sufficient for self-adjointness.

Quantum Graphs: Finite total volume

$\text{vol}(\mathcal{G}) = \sum_{e \in \mathcal{E}} |e| < \infty$, then \mathbf{H} is non-self-adjoint

Problem

Deficiency indices? Self-adjoint extensions? Boundary conditions?

Quantum Graphs: Finite total volume

$\text{vol}(\mathcal{G}) = \sum_{e \in \mathcal{E}} |e| < \infty$, then \mathbf{H} is non-self-adjoint

Problem

Deficiency indices? Self-adjoint extensions? Boundary conditions?

Graph Ends

- A ray R in \mathcal{G}_d is a path without intersections.

Quantum Graphs: Finite total volume

$\text{vol}(\mathcal{G}) = \sum_{e \in \mathcal{E}} |e| < \infty$, then \mathbf{H} is non-self-adjoint

Problem

Deficiency indices? Self-adjoint extensions? Boundary conditions?

Graph Ends

- A ray R in \mathcal{G}_d is a path without intersections.
- Two rays are *equivalent* if there is a third ray containing infinitely many vertices of both rays.

Quantum Graphs: Finite total volume

$\text{vol}(\mathcal{G}) = \sum_{e \in \mathcal{E}} |e| < \infty$, then \mathbf{H} is non-self-adjoint

Problem

Deficiency indices? Self-adjoint extensions? Boundary conditions?

Graph Ends

- A ray R in \mathcal{G}_d is a path without intersections.
- Two rays are *equivalent* if there is a third ray containing infinitely many vertices of both rays.
- An equivalence class of rays is a *graph end*; $\Omega(\mathcal{G}_d)$ is the set of graph ends.

Theorem (e.g., Diestel–Kühn '2003)

Topological ends of \mathcal{G} = graph ends of \mathcal{G}_d .

Graph Ends: Examples

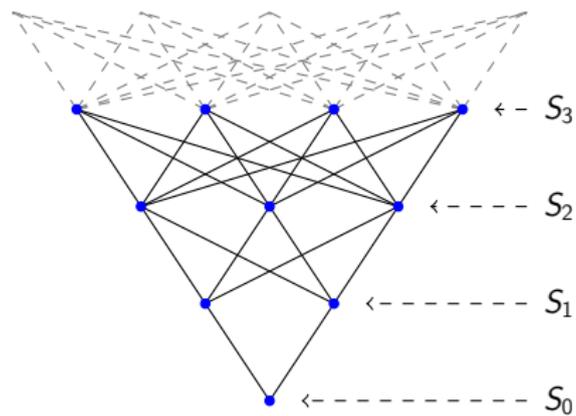


Figure: An antitree \mathcal{A} with $s_n = n + 1$.

Every antitree has exactly 1 end.

Graph Ends: Examples

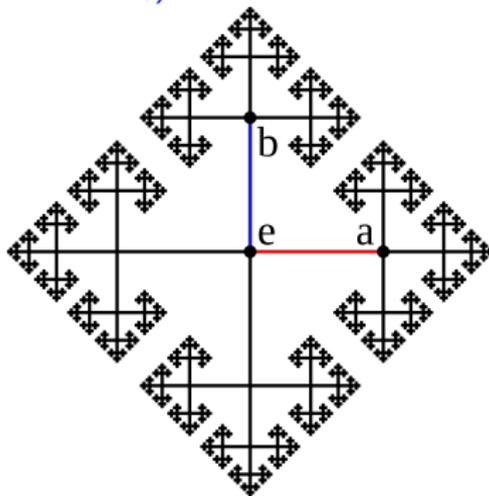
- $\mathcal{G}_d = \mathbb{Z}$ has 2 ends.

Graph Ends: Examples

- $\mathcal{G}_d = \mathbb{Z}$ has 2 ends.
- $\mathcal{G}_d = \mathbb{Z}^N$ has 1 end for all $N \geq 2$.

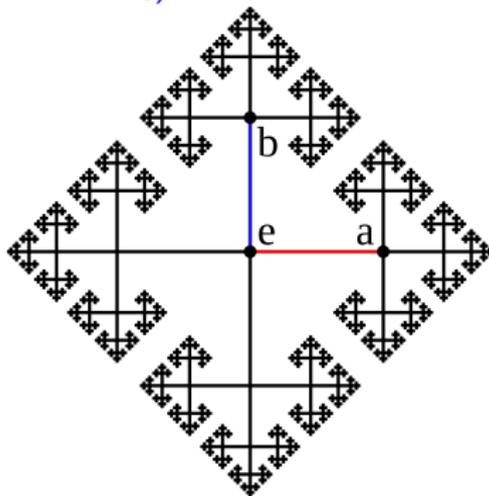
Graph Ends: Examples

- $\mathcal{G}_d = \mathbb{Z}$ has 2 ends.
- $\mathcal{G}_d = \mathbb{Z}^N$ has 1 end for all $N \geq 2$.
- Bethe lattice (Cayley or regular tree \mathbb{T}_4)



Graph Ends: Examples

- $\mathcal{G}_d = \mathbb{Z}$ has 2 ends.
- $\mathcal{G}_d = \mathbb{Z}^N$ has 1 end for all $N \geq 2$.
- Bethe lattice (Cayley or regular tree \mathbb{T}_4)



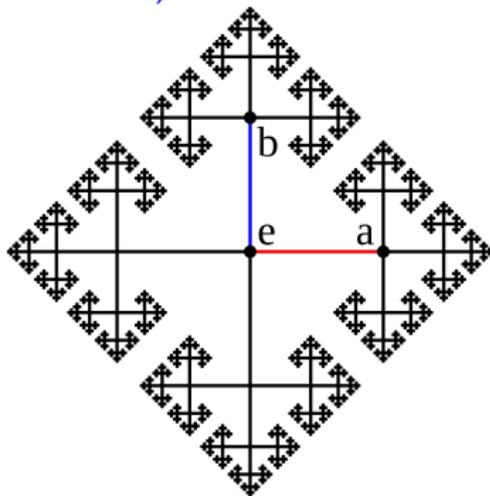
Theorem (J. R. Stallings, *Ann. of Math.* (1968))

If \mathcal{G}_d is a Cayley graph of a finitely generated group, then

$\#\Omega(\mathcal{G}_d)$

Graph Ends: Examples

- $\mathcal{G}_d = \mathbb{Z}$ has 2 ends.
- $\mathcal{G}_d = \mathbb{Z}^N$ has 1 end for all $N \geq 2$.
- Bethe lattice (Cayley or regular tree \mathbb{T}_4)



Theorem (J. R. Stallings, *Ann. of Math.* (1968))

If \mathcal{G}_d is a Cayley graph of a finitely generated group, then $\#\Omega(\mathcal{G}_d) \in \{1, 2, \infty\}$.

Theorem (AK–Mugnolo–Nicolussi, *in preparation*)

If $\text{vol}(\mathcal{G}) < \infty$, then $n_{\pm}(\mathbf{H}) \geq \#\Omega(\mathcal{G}_d)$.

Quantum Graphs: Deficiency Indices

Theorem (AK–Mugnolo–Nicolussi, *in preparation*)

If $\text{vol}(\mathcal{G}) < \infty$, then $n_{\pm}(\mathbf{H}) \geq \#\Omega(\mathcal{G}_d)$.

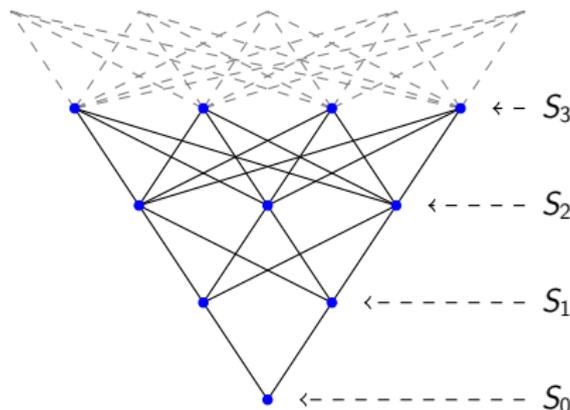


Figure: An antitree \mathcal{A} with $s_n = n + 1$.

For radially symmetric antitrees, $n_{\pm}(\mathcal{A}) = 1$ iff $\text{vol}(\mathcal{A}) < \infty$

Quantum Graphs: Deficiency Indices

Theorem (AK–Mugnolo–Nicolussi, *in preparation*)

If $\text{vol}(\mathcal{G}) < \infty$, then $n_{\pm}(\mathbf{H}) \geq \#\Omega(\mathcal{G}_d)$.

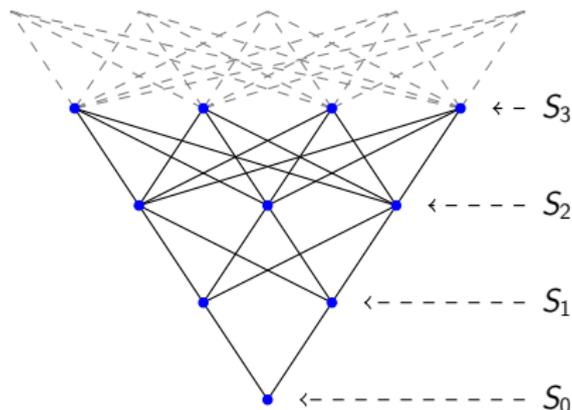


Figure: An antitree \mathcal{A} with $s_n = n + 1$.

For radially symmetric antitrees, $n_{\pm}(\mathcal{A}) = 1$ iff $\text{vol}(\mathcal{A}) < \infty$

However, there are antitrees with $n_{\pm}(\mathcal{A}) = \infty$!

Theorem (AK–Mugnolo–Nicolussi, *in preparation*)

If $\text{vol}(\mathcal{G}) < \infty$, then $n_{\pm}(\mathbf{H}) \geq \#\Omega(\mathcal{G}_d)$. Moreover,

$n_{\pm}(\mathbf{H}) = \#\Omega(\mathcal{G}_d)$ if and only if either $\#\Omega(\mathcal{G}_d) = \infty$ or $\ker(\mathbf{H}^*) \subset H^1(\mathcal{G})$.

Here $H^1(\mathcal{G})$ is the usual Sobolev space on \mathcal{G} .

Theorem (AK–Mugnolo–Nicolussi, *in preparation*)

If $\text{vol}(\mathcal{G}) < \infty$, then $n_{\pm}(\mathbf{H}) \geq \#\Omega(\mathcal{G}_d)$. Moreover,

$n_{\pm}(\mathbf{H}) = \#\Omega(\mathcal{G}_d)$ if and only if either $\#\Omega(\mathcal{G}_d) = \infty$ or $\ker(\mathbf{H}^*) \subset H^1(\mathcal{G})$.

Here $H^1(\mathcal{G})$ is the usual Sobolev space on \mathcal{G} .

Remarks

- Since 0 is a point of a regular type for \mathbf{H} , $n_{\pm}(\mathbf{H}) = \dim(\ker(\mathbf{H}^*))$.

Theorem (AK–Mugnolo–Nicolussi, *in preparation*)

If $\text{vol}(\mathcal{G}) < \infty$, then $n_{\pm}(\mathbf{H}) \geq \#\Omega(\mathcal{G}_d)$. Moreover, $n_{\pm}(\mathbf{H}) = \#\Omega(\mathcal{G}_d)$ if and only if either $\#\Omega(\mathcal{G}_d) = \infty$ or $\ker(\mathbf{H}^*) \subset H^1(\mathcal{G})$.

Here $H^1(\mathcal{G})$ is the usual Sobolev space on \mathcal{G} .

Remarks

- Since 0 is a point of a regular type for \mathbf{H} , $n_{\pm}(\mathbf{H}) = \dim(\ker(\mathbf{H}^*))$.
- $\ker(\mathbf{H}^*)$ consists of harmonic functions which belong to $L^2(\mathcal{G})$.

Theorem (AK–Mugnolo–Nicolussi, *in preparation*)

If $\text{vol}(\mathcal{G}) < \infty$, then $n_{\pm}(\mathbf{H}) \geq \#\Omega(\mathcal{G}_d)$. Moreover,

$n_{\pm}(\mathbf{H}) = \#\Omega(\mathcal{G}_d)$ if and only if either $\#\Omega(\mathcal{G}_d) = \infty$ or $\ker(\mathbf{H}^*) \subset H^1(\mathcal{G})$.

Here $H^1(\mathcal{G})$ is the usual Sobolev space on \mathcal{G} .

Remarks

- Since 0 is a point of a regular type for \mathbf{H} , $n_{\pm}(\mathbf{H}) = \dim(\ker(\mathbf{H}^*))$.
- $\ker(\mathbf{H}^*)$ consists of harmonic functions which belong to $L^2(\mathcal{G})$.
- $H^1(\mathcal{G})$ is a 'nice' space (e.g., graph ends can be identified with its Royden's boundary, which gives a hope for reasonable traces of functions in $\text{dom}(\mathbf{H}^*)$).

In the discrete setting, see



A. Georgakopoulos, S. Haeseler, M. Keller, D. Lenz and R. Wojciechowski, *Graphs of finite measure*, J. Math. Pures Appl. **103** (2015).

Weighted Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$.

Suppose we are given two more edge weights

$$\mu: \mathcal{E} \rightarrow \mathbb{R}_{>0},$$

$$\nu: \mathcal{E} \rightarrow \mathbb{R}_{>0}$$

Weighted Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$.

Suppose we are given two more edge weights

$$\mu: \mathcal{E} \rightarrow \mathbb{R}_{>0},$$

$$\nu: \mathcal{E} \rightarrow \mathbb{R}_{>0}$$

Introduce the *weighted* Hilbert space $L^2(\mathcal{G}; \mu) := \bigoplus_{e \in \mathcal{E}} L^2(e; \mu_e)$ and equip \mathcal{G} with a Schrödinger-type operator $\mathbf{H}_{\max} := \bigoplus_{e \in \mathcal{E}} H_{\mu, \nu}^e$, where:

$$H_{\mu, \nu}^e = -\frac{1}{\mu_e} \frac{d}{dx_e} \nu_e \frac{d}{dx_e}, \quad \text{dom}(H_{\mu, \nu}^e) = H^2(e).$$

Weighted Quantum Graphs

Given a metric graph $\mathcal{G} = (\mathcal{V}, \mathcal{E}, |\cdot|)$.

Suppose we are given two more edge weights

$$\mu: \mathcal{E} \rightarrow \mathbb{R}_{>0}, \quad \nu: \mathcal{E} \rightarrow \mathbb{R}_{>0}$$

Introduce the *weighted* Hilbert space $L^2(\mathcal{G}; \mu) := \bigoplus_{e \in \mathcal{E}} L^2(e; \mu_e)$ and equip \mathcal{G} with a Schrödinger-type operator $\mathbf{H}_{\max} := \bigoplus_{e \in \mathcal{E}} H_{\mu, \nu}^e$, where:

$$H_{\mu, \nu}^e = -\frac{1}{\mu_e} \frac{d}{dx_e} \nu_e \frac{d}{dx_e}, \quad \text{dom}(H_{\mu, \nu}^e) = H^2(e).$$

The operator $\mathbf{H}_{\mu, \nu}$ with Kirchhoff conditions: For all $v \in \mathcal{V}$

$$\begin{cases} f \text{ is continuous at } v, \\ \sum_{e \in \mathcal{E}_v} \nu_e f'_e(v) = 0. \end{cases}$$

Weighted Quantum Graphs

The analog of Theorem 1 for $\mathbf{H}_{\mu,\nu}$ holds true, however, with the (minimal) discrete Laplacian defined on $\ell^2(\mathcal{V}; m_\mu)$ by

$$(\tau_G f)(v) := \frac{1}{m_\mu(v)} \sum_{u \sim v} b_\nu(e_{u,v})(f(v) - f(u)),$$

where

$$m_\mu(v) = \sum_{e \in \mathcal{E}_v} \mu_e |e|, \quad b_\nu(e) = \frac{\nu_e}{|e|}.$$

Weighted Quantum Graphs

The analog of Theorem 1 for $\mathbf{H}_{\mu,\nu}$ holds true, however, with the (minimal) discrete Laplacian defined on $\ell^2(\mathcal{V}; m_\mu)$ by

$$(\tau_{\mathcal{G}}f)(v) := \frac{1}{m_\mu(v)} \sum_{u \sim v} b_\nu(e_{u,v})(f(v) - f(u)),$$

where

$$m_\mu(v) = \sum_{e \in \mathcal{E}_v} \mu_e |e|, \quad b_\nu(e) = \frac{\nu_e}{|e|}.$$

Remark (self-adjointness)

If \mathcal{G} is a path graph, then $\mathbf{H}_{\mu,\nu}$ is self-adjoint if and only if

$$\sum_n \mu_n |e_n| \left(\sum_{k \leq n} \frac{|e_k|}{\nu_k} \right)^2 = \infty.$$

Weighted Quantum Graphs

The analog of Theorem 1 for $\mathbf{H}_{\mu,\nu}$ holds true, however, with the (minimal) discrete Laplacian defined on $\ell^2(\mathcal{V}; m_\mu)$ by

$$(\tau_{\mathcal{G}}f)(v) := \frac{1}{m_\mu(v)} \sum_{u \sim v} b_\nu(e_{u,v})(f(v) - f(u)),$$

where

$$m_\mu(v) = \sum_{e \in \mathcal{E}_v} \mu_e |e|, \quad b_\nu(e) = \frac{\nu_e}{|e|}.$$

Remark (self-adjointness)

If \mathcal{G} is a path graph, then $\mathbf{H}_{\mu,\nu}$ is self-adjoint if and only if

$$\sum_n \mu_n |e_n| \left(\sum_{k \leq n} \frac{|e_k|}{\nu_k} \right)^2 = \infty.$$

Hence $\sum m_\mu(v_n) = 2 \sum \mu_n |e_n| = \infty$ is only sufficient!

Weighted Quantum Graphs

The analog of Theorem 1 for $\mathbf{H}_{\mu,\nu}$ holds true, however, with the (minimal) discrete Laplacian defined on $\ell^2(\mathcal{V}; m_\mu)$ by

$$(\tau_G f)(v) := \frac{1}{m_\mu(v)} \sum_{u \sim v} b_\nu(e_{u,v})(f(v) - f(u)),$$

where

$$m_\mu(v) = \sum_{e \in \mathcal{E}_v} \mu_e |e|, \quad b_\nu(e) = \frac{\nu_e}{|e|}.$$

Weighted discrete Laplacian

For $m: \mathcal{V} \rightarrow \mathbb{R}_{>0}$ and $b: \mathcal{E} \rightarrow \mathbb{R}_{>0}$, consider in $\ell^2(\mathcal{V}; m)$

$$(\tau f)(v) := \frac{1}{m(v)} \sum_{u \sim v} b(e_{u,v})(f(v) - f(u)).$$

Weighted Quantum Graphs

The analog of Theorem 1 for $\mathbf{H}_{\mu,\nu}$ holds true, however, with the (minimal) discrete Laplacian defined on $\ell^2(\mathcal{V}; m_\mu)$ by

$$(\tau_{\mathcal{G}}f)(v) := \frac{1}{m_\mu(v)} \sum_{u \sim v} b_\nu(e_{u,v})(f(v) - f(u)),$$

where

$$m_\mu(v) = \sum_{e \in \mathcal{E}_v} \mu_e |e|, \quad b_\nu(e) = \frac{\nu_e}{|e|}.$$

Weighted discrete Laplacian

For $m: \mathcal{V} \rightarrow \mathbb{R}_{>0}$ and $b: \mathcal{E} \rightarrow \mathbb{R}_{>0}$, consider in $\ell^2(\mathcal{V}; m)$

$$(\tau f)(v) := \frac{1}{m(v)} \sum_{u \sim v} b(e_{u,v})(f(v) - f(u)).$$

QUESTION: For a given τ (i.e., a pair of functions m and b), does there exist a “weighted” \mathcal{G} (i.e., weights $|\cdot|$, μ and ν) such that $\tau = \tau_{\mathcal{G}}$?

Normalized/Physical Laplacian

Take $\mu_e = \nu_e = |e|$ for all $e \in \mathcal{E}$, then

$$m_\mu(v) = \deg(v), \quad b_\nu(e) = 1.$$

Weighted Quantum Graphs: Examples

Normalized/Physical Laplacian

Take $\mu_e = \nu_e = |e|$ for all $e \in \mathcal{E}$, then

$$m_\mu(v) = \deg(v), \quad b_\nu(e) = 1.$$

Electric Networks/Random Walks on Graphs

Take $\nu_e = |e|b(e)$ and $\mu_e = \frac{b(e)}{|e|}$ for all $e \in \mathcal{E}$, then

$$m_\mu(v) = \sum_{e \in \mathcal{E}} b(e) = m(e), \quad b_\nu(e) = b(e).$$

Weighted Quantum Graphs: Examples

Normalized/Physical Laplacian

Take $\mu_e = \nu_e = |e|$ for all $e \in \mathcal{E}$, then

$$m_\mu(v) = \deg(v), \quad b_\nu(e) = 1.$$

Electric Networks/Random Walks on Graphs

Take $\nu_e = |e|b(e)$ and $\mu_e = \frac{b(e)}{|e|}$ for all $e \in \mathcal{E}$, then

$$m_\mu(v) = \sum_{e \in \mathcal{E}} b(e) = m(e), \quad b_\nu(e) = b(e).$$

Path Graphs and Jacobi Matrices

Weighted Quantum Graphs: Examples

Normalized/Physical Laplacian

Take $\mu_e = \nu_e = |e|$ for all $e \in \mathcal{E}$, then

$$m_\mu(v) = \deg(v), \quad b_\nu(e) = 1.$$

Electric Networks/Random Walks on Graphs

Take $\nu_e = |e|b(e)$ and $\mu_e = \frac{b(e)}{|e|}$ for all $e \in \mathcal{E}$, then

$$m_\mu(v) = \sum_{e \in \mathcal{E}} b(e) = m(e), \quad b_\nu(e) = b(e).$$

Path Graphs and Jacobi Matrices

Every Jacobi matrix can be realized as a boundary operator for a weighted quantum path graph (with δ -interactions at the vertices)

Weighted Quantum Graphs

Combinatorial Laplacian: $m \equiv 1, b \equiv 1$

$$(\tau_{\text{comb}} f)(v) := \sum_{u \sim v} f(v) - f(u) = \deg(v)f(v) - \underbrace{\sum_{u \sim v} f(u)}_{\text{adjacency matrix}} .$$

Weighted Quantum Graphs

Combinatorial Laplacian: $m \equiv 1, b \equiv 1$

$$(\tau_{\text{comb}} f)(v) := \sum_{u \sim v} f(v) - f(u) = \deg(v)f(v) - \underbrace{\sum_{u \sim v} f(u)}_{\text{adjacency matrix}} .$$

If \mathcal{G}_d has loose ends, then one can't construct a metric with $\mathbf{H}_{\mu,\nu}$!

Weighted Quantum Graphs

Combinatorial Laplacian: $m \equiv 1, b \equiv 1$

$$(\tau_{\text{comb}} f)(v) := \sum_{u \sim v} f(v) - f(u) = \deg(v)f(v) - \underbrace{\sum_{u \sim v} f(u)}_{\text{adjacency matrix}} .$$

If \mathcal{G}_d has loose ends, then one can't construct a metric with $\mathbf{H}_{\mu,\nu}$!

For an antitree \mathcal{A} , only if $\sum_{k=0}^n (-1)^k s_{n-k} > 0$ for all $n \geq 0$.

Weighted Quantum Graphs

Combinatorial Laplacian: $m \equiv 1$, $b \equiv 1$

$$(\tau_{\text{comb}} f)(v) := \sum_{u \sim v} f(v) - f(u) = \deg(v)f(v) - \underbrace{\sum_{u \sim v} f(u)}_{\text{adjacency matrix}} .$$

If \mathcal{G}_d has loose ends, then one can't construct a metric with $\mathbf{H}_{\mu,\nu}$!

For an antitree \mathcal{A} , only if $\sum_{k=0}^n (-1)^k s_{n-k} > 0$ for all $n \geq 0$.

Theorem (G. Zaimi '2011: mathoverflow.net/questions/59117)

Let $\mathcal{G}_d = (\mathcal{V}, \mathcal{E})$ be simple, connected, locally finite. Then there are lengths $|\cdot|: \mathcal{E} \rightarrow \mathbb{R}_{>0}$ and weights $\mu: \mathcal{E} \rightarrow \mathbb{R}_{>0}$ such that

$$\sum_{e \in \mathcal{E}_v} \mu(e)|e| = 1 \quad \text{for all } v \in \mathcal{V},$$

if and only if for each $e \in \mathcal{E}$ there is a disjoint cycle cover containing e in one of its cycles.

Weighted Quantum Graphs

Combinatorial Laplacian: $m \equiv 1$, $b \equiv 1$

$$(\tau_{\text{comb}} f)(v) := \sum_{u \sim v} f(v) - f(u) = \deg(v)f(v) - \underbrace{\sum_{u \sim v} f(u)}_{\text{adjacency matrix}} .$$

If \mathcal{G}_d has loose ends, then one can't construct a metric with $\mathbf{H}_{\mu,\nu}$!

For an antitree \mathcal{A} , only if $\sum_{k=0}^n (-1)^k s_{n-k} > 0$ for all $n \geq 0$.

The way to fix this problem is to **allow loops!**



M. Folz, *Volume growth and stochastic completeness of graphs*, Trans. Amer. Math. Soc. **366** (2014).



X. Huang, *A note on the volume growth criterion for stochastic completeness of weighted graphs*, Potential Anal. **40** (2014).

Weighted Quantum Graphs

Combinatorial Laplacian: $m \equiv 1, b \equiv 1$

$$(\tau_{\text{comb}} f)(v) := \sum_{u \sim v} f(v) - f(u) = \deg(v)f(v) - \underbrace{\sum_{u \sim v} f(u)}_{\text{adjacency matrix}} .$$

If \mathcal{G}_d has loose ends, then one can't construct a metric with $\mathbf{H}_{\mu,\nu}$!

For an antitree \mathcal{A} , only if $\sum_{k=0}^n (-1)^k s_{n-k} > 0$ for all $n \geq 0$.

The way to fix this problem is to **allow loops!**

Then **every weighted discrete Laplacian** can be realized as a **boundary operator for a quantum graph** operator (in the sense of Theorem 1), however, the metric graph might be with loops.



A. Kostenko, M. Malamud, and N. Nicolussi, *Weighted quantum graphs*, in preparation.

8th ECM in Portorož, Slovenia: July 5–11, 2020



8th ECM in Portorož, Slovenia: July 5–11, 2020



Thank you for your attention!