# Magnetic Schrödinger operators with electric $\delta$-potentials 

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Graz University of Technology
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## Outline

1. Motivation
2. Magnetic Schrödinger operators with $\delta$-potentials

- The magnetic Schrödinger operator without potential
- Magnetic Sobolev spaces
- Definition of the $\delta$-operator

3. Approximation by Hamiltonians with squeezed potentials
4. Exner-Ichinose for homogeneous magnetic fields
5. A quasi boundary triple
6. Outlook

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## Schrödinger operator with magnetic fields

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\left(i \partial_{t}-\left(-i \nabla_{x}-A\right)^{2}+V\right) \psi(t, x)=0 \quad+\quad \text { i. c. }
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where $B=\nabla \times A$, i. e. $A: \mathbb{R}^{3} \rightarrow \mathbb{R}^{3}$

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- We consider $H$ in $L^{2}\left(\mathbb{R}^{d}\right)$ for any $d \geq 2$ (physical meaning for $d=2,3$ )


## Hamiltonians with $\delta$-potentials

For a zero-set $\Sigma \subset \mathbb{R}^{d}$ and $\alpha: \Sigma \rightarrow \mathbb{R}$ consider

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## Conjectures:

- For homogeneous magnetic fields ( $B=$ const.): same behavior
- For non-homogeneous fields: bound states disappear


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- $\mathfrak{h}_{0}$ is densely defined, closed, and $\mathfrak{h}_{0} \geq 0$
- associated self-adjoint operator

$$
H_{0}:=(-i \nabla-A)^{2}
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## Definition of magnetic Sobolev spaces

- Problem: for $A \in C^{\infty}\left(\mathbb{R}^{d} ; \mathbb{R}^{d}\right)$ we have in general $f \in H^{1}\left(\mathbb{R}^{d}\right) \nRightarrow f \in \mathcal{H}_{A}^{1}\left(\mathbb{R}^{d}\right)$


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- $\mathcal{H}_{A}^{s}(\Omega)$, equipped with the natural norm, is a Hilbert space


## The diamagnetic inequality

## Theorem

Let $t>0$ and $f \in L^{2}\left(\mathbb{R}^{d}\right)$. Then:

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\left|e^{-t H_{0}} f\right| \leq e^{-t(-\Delta)}|f| .
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## Consequences:

- It holds for $s>0, r \geq 0$, and $\lambda<0$

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(r-\lambda)^{-s}=\frac{1}{\Gamma(-\lambda)} \int_{0}^{\infty} t^{s-1} e^{-t(r-\lambda)} \mathrm{d} t
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## Corollary

$\mathcal{H}_{A}^{s}\left(\mathbb{R}^{d}\right) \subset H^{s}\left(\mathbb{R}^{d}\right)$ for all $s \geq 0$.

## Definition of the $\delta$-operator - preparations

- Let $\left\{\Sigma_{j}\right\}_{j=1}^{N}$ be a family of smooth hypersurfaces with $\sigma\left(\Sigma_{k} \cap \Sigma_{l}\right)=0, k \neq 1$


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- Since $\mathcal{H}_{A}^{1}\left(\mathbb{R}^{d}\right) \subset H^{1}\left(\mathbb{R}^{d}\right)$, the trace $\left.f\right|_{\Sigma} \in L^{2}(\Sigma)$ for $f \in \mathcal{H}_{A}^{1}\left(\mathbb{R}^{d}\right)$


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& \text { Proof: } \forall a>0 \exists b>0: \\
& \qquad \mathfrak{h}_{\Sigma}[f] \leq a\|\nabla f\|^{2}+b\|f\|^{2}
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Proof: $\forall a>0 \exists b>0$ :

$$
\mathfrak{h}_{\Sigma}[f] \leq a\|\nabla f\|^{2}+b\|f\|^{2} \leq a \mathfrak{h}_{0}[f]+b\|f\|^{2}
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(diamagnetic inequality)

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- Remark: One can add a form bounded potential $Q$ with relative bound < 1


## Outline

1. Motivation
2. Magnetic Schrödinger operators with $\delta$-potentials

- The magnetic Schrödinger operator without potential
- Magnetic Sobolev spaces
- Definition of the $\delta$-operator

3. Approximation by Hamiltonians with squeezed potentials
4. Exner-Ichinose for homogeneous magnetic fields
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## Construction of the approximating sequence

- Assume $\exists \beta>0$ such that
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- $(-i \nabla-A)^{2}-\sum_{j=1}^{N} V_{j, \varepsilon}$ is self-adjoint on $\mathcal{H}_{A}^{2}\left(\mathbb{R}^{d}\right)$


## The result

## Theorem

Define $\alpha \in L^{\infty}(\Sigma)$ as

$$
\alpha\left(x_{\Sigma}\right):=\int_{-\beta}^{\beta} V_{j}\left(x_{\Sigma}+\boldsymbol{s} \nu\left(x_{\Sigma}\right)\right) \mathrm{d} s, \quad x_{\Sigma} \in \Sigma_{j}
$$

and let $\lambda \ll 0$. Then there exists $c>0$ such that

$$
\left\|\left((-i \nabla-A)^{2}-\sum_{j=1}^{N} V_{j, \varepsilon}-\lambda\right)^{-1}-\left(H_{\alpha}-\lambda\right)^{-1}\right\| \leq c \varepsilon
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for small $\varepsilon>0$. In particular $(-i \nabla-A)^{2}-\sum_{j=1}^{N} V_{j, \varepsilon}$ converge to $H_{\alpha}$ in the norm resolvent sense.

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- Adding a form bounded potential $Q$ does not change the argument


## Comparison to [Behrndt-Exner-H-Lotoreichik'17]

Known
New

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## Statement

Magnetic field

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- Physical interpretation: $(0,0, B)^{\top}=\nabla \times(A, 0)^{\top}$, i.e. the magnetic field is perpendicular to the plane
- For $\lambda \in \rho\left((-i \nabla-A)^{2}\right)$ it holds

$$
\left((-i \nabla-A)^{2}-\lambda\right)^{-1} f(x)=\int_{\mathbb{R}^{2}} G_{\lambda}^{A}(x, y) f(y) \mathrm{d} y
$$

where $G_{\lambda}^{A}$ is explicitely given by a combination of

- an irregular confluent hypergeometric function
- an in general complex valued function


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Assume that $\Sigma \neq \Gamma$ and $\alpha>0$ is constant. Then, $\sigma_{\text {disc }}\left(-\Delta-\alpha \delta_{\Sigma}\right) \neq \emptyset$.

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Goal: Prove a similar result for $B \neq 0$

## Proof of Exner-Ichinose

- Use Birman-Schwinger principle:

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\begin{aligned}
&(-\infty, 0) \ni \lambda \in \sigma\left(-\Delta-\alpha \delta_{\Sigma}\right) \Leftrightarrow 1 \in \sigma\left(\alpha M_{\Sigma}(\lambda)\right) \\
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## Outline

## 1. Motivation

2. Magnetic Schrödinger operators with $\delta$-potentials

- The magnetic Schrödinger operator without potential
- Magnetic Sobolev spaces
- Definition of the $\delta$-operator

3. Approximation by Hamiltonians with squeezed potentials
4. Exner-Ichinose for homogeneous magnetic fields
5. A quasi boundary triple
6. Outlook

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S f:=(-i \nabla-A)^{2} f, \quad \operatorname{dom} S=\left\{f \in \mathcal{H}_{A}^{2}\left(\mathbb{R}^{d}\right):\left.f\right|_{\Sigma}=0\right\}
$$

## Some notations

- In this section: $\Sigma \subset \mathbb{R}^{d}$ is the boundary of a smooth bounded domain $\Omega_{+}$with outer normal vector field $\nu$
- $\Omega_{-}:=\mathbb{R}^{d} \backslash \overline{\Omega_{+}}$
- Define

$$
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and

$$
\begin{aligned}
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\end{aligned}
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## The quasi boundary triple

$$
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- It holds for all $f, g \in \operatorname{dom} T$

$$
\begin{aligned}
& \left((-i \nabla-A)^{2} f_{+}, g_{+}\right)_{\Omega_{+}}-\left(f_{+},(-i \nabla-A)^{2} g_{+}\right)_{\Omega_{+}} \\
& \quad=\left(\left.f\right|_{\Sigma},\left(\partial_{\nu}-i \nu A\right) g_{+} \mid \Sigma\right)_{\Sigma}-\left(\left.\left(\partial_{\nu}-i \nu A\right) f_{+}\right|_{\Sigma},\left.g\right|_{\Sigma}\right)_{\Sigma}
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and hence

$$
(T f, g)_{\mathbb{R}^{d}}-(f, T g)_{\mathbb{R}^{d}}=\left(\Gamma_{1} f, \Gamma_{0} g\right)_{\Sigma}-\left(\Gamma_{0} f, \Gamma_{1} g\right)_{\Sigma}
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- $\operatorname{ran}\left(\Gamma_{0}, \Gamma_{1}\right)=H^{1 / 2}(\Sigma) \times H^{3 / 2}(\Sigma)$
- $A_{0}:=T \upharpoonright \operatorname{ker} \Gamma_{0}$ is the free operator $(-i \nabla-A)^{2}$ in $L^{2}\left(\mathbb{R}^{d}\right)$


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- $\Rightarrow\left\{L^{2}(\Sigma), \Gamma_{0}, \Gamma_{1}\right\}$ is a quasi boundary triple for $S^{*}$


## The quasi boundary triple and the $\delta$-operator

Define for $\alpha \in \mathbb{R}$ the operator $H_{\alpha}^{Q}:=T \upharpoonright \operatorname{ker}\left(\Gamma_{0}-\alpha \Gamma_{1}\right)$, i.e.

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## Theorem

$A_{\alpha}^{Q}$ is self-adjoint and coincides with $A_{\alpha}$. In particular, $\operatorname{dom} A_{\alpha} \subset \mathcal{H}_{A}^{2}\left(\Omega_{+}\right) \oplus \mathcal{H}_{A}^{2}\left(\Omega_{-}\right)$.

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- Green's formula: $A_{\alpha}^{Q}$ is symmetric


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- Formulae for scattering theory


## Outline

## 1. Motivation

2. Magnetic Schrödinger operators with $\delta$-potentials

- The magnetic Schrödinger operator without potential
- Magnetic Sobolev spaces
- Definition of the $\delta$-operator

3. Approximation by Hamiltonians with squeezed potentials
4. Exner-Ichinose for homogeneous magnetic fields
5. A quasi boundary triple
6. Outlook

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## Thank you for your attention!

